

Riemannian Lie algebroids

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Seminar Algebra, Geometry, Topology and Applications

Outline

- 1 Connections on Lie algebroids
 - Parallel transport
 - Linear A -connections, geodesics and compatibility with the Lie algebroid structure
 - Variations of A -paths, homotopy and curvature of A -connections
 - Homotopy of A -paths
- 2 Riemannian metrics on Lie algebroids
- 3 Geodesic flow of a Riemannian Lie algebroid
- 4 First and second variation formula
- 5 O'Neill's formulas for curvature
- 6 Integrability of Riemannian Lie algebroids

An *anchored vector bundle* is a triple (A, M, ρ) where :

- 1 $p : A \rightarrow M$ is a vector bundle,
- 2 $\rho : A \rightarrow TM$ is a bundle homomorphism called anchor.

Let (A, M, ρ) be an anchored vector bundle.

A bracket on $\Gamma(A)$ is a skew-symmetric \mathbb{R} -bilinear map

$$[\cdot, \cdot]_A : \Gamma(A) \times \Gamma(A) \rightarrow \Gamma(A).$$

It is called *anchored* if for any $a, b \in \Gamma(A)$ and for every smooth function $f \in C^\infty(M)$,

$$[a, fb]_A = f[a, b]_A + \rho(a)(f)b. \quad (1)$$

$[\cdot, \cdot]_A$ is local:

$$a|_U = 0 \implies (\forall b \in \Gamma(A), \quad [a, b]_A = 0 \text{ on } U.)$$

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The *torsion* of $[\ , \]_A$ is the map $\Theta_{[\ , \]_A} : \Gamma(A) \times \Gamma(A) \longrightarrow \mathcal{X}(M)$ given by

$$\Theta_{[\ , \]_A}(a, b) = \rho([a, b]_A) - [\rho(a), \rho(b)]. \quad (2)$$

$\Theta_{[\ , \]_A}$ is \mathbb{R} -bilinear, skew-symmetric and, for any $f \in C^\infty(M)$,

$$\Theta_{[\ , \]_A}(fa, b) = \Theta_{[\ , \]_A}(a, fb) = f\Theta_{[\ , \]_A}(a, b).$$

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The Jacobiator of $[\cdot, \cdot]_A$ as $J_{[\cdot, \cdot]_A} : \Gamma(A) \times \Gamma(A) \times \Gamma(A) \longrightarrow \Gamma(A)$ given by

$$J_{[\cdot, \cdot]_A}(a, b, c) = [[a, b]_A, c]_A + [[b, c]_A, a]_A + [[c, a]_A, b]_A.$$

J is \mathbb{R} -trilinear and skew-symmetric. Thus $[\cdot, \cdot]_A$ is a Lie bracket if and only if $J_{[\cdot, \cdot]_A} = 0$.

For any $a, b, c \in \Gamma(A)$ and any $f \in C^\infty(M)$,

$$J_{[\cdot, \cdot]_A}(a, b, fc) = fJ_{[\cdot, \cdot]_A}(a, b, c) + \Theta_{[\cdot, \cdot]_A}(a, b)(f)c. \quad (3)$$

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This relation shows that:

- 1 $J_{[\cdot, \cdot]_A}$ is local:

$$a|_U = 0 \implies J_A(a, \cdot, \cdot)|_U = 0,$$

- 2 if $\Theta_{[\cdot, \cdot]_A}$ vanishes then $J_{[\cdot, \cdot]_A}$ becomes a tensor, namely,
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Proposition

Let (A, M, ρ) be an anchored bundle and $[\cdot, \cdot]_A$ an anchored bracket on $\Gamma(A)$. Then the following assertions are equivalent:

- (i) $(\Gamma(A), [\cdot, \cdot]_A)$ is a Lie algebra, i.e., J_A vanishes identically.
- (ii) For any $x \in M$ there exists an open set U of M containing x and a basis of sections (a_1, \dots, a_r) over U such that

$$J_{[\cdot, \cdot]_A}(a_i, a_j, a_k) = 0 \quad \text{and} \quad \Theta(a_i, a_j) = 0, \quad 1 \leq i < j < k \leq r.$$

Definition

A Lie algebroid is an anchored vector bundle (A, M, ρ) together with an anchored bracket $[\cdot, \cdot]_A$ satisfying (i) or (ii) of Proposition 1.3.

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There are some well-known properties of a Lie algebroid $(A, M, \rho, [\cdot, \cdot]_A)$.

- (a) The smooth distribution $\text{Im}\rho$ is integrable in the sense of Sussmann and, for any leaf L of $\text{Im}\rho$, $(A|_L, L, \rho, [\cdot, \cdot]_A)$ is a transitive Lie algebroid.
- (b) For any $x \in M$, there is an induced Lie bracket say $[\cdot, \cdot]_x$ on $\mathfrak{g}_x = \ker(\rho_x) \subset A_x$ which makes it into a finite dimensional Lie algebra.

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(c) The map $d_A : \Gamma(\wedge A^*) \longrightarrow \Gamma(\wedge A^*)$ by

$$d_A Q(a_1, \dots, a_p) = \sum_{i=1}^p (-1)^{i+1} \rho(a_i) \cdot Q(a_1, \dots, \hat{a}_i, \dots, a_p) \\ - \sum_{1 \leq i < j \leq p} (-1)^{i+j+1} Q([a_i, a_j]_A, a_1, \dots, \hat{a}_i, \dots, \hat{a}_j, \dots, a_p),$$

is a differential, i.e., $d_A^2 = 0$. In particular, for any $a, b \in \Gamma(A)$, $f \in C^\infty(M)$ and $Q \in \Gamma(\wedge A^*)$,

$$d_A f(a) = \rho(a)(f)$$

$$d_A Q(a, b) = \rho(a) \cdot Q(b) - \rho(b) \cdot Q(a) - Q([a, b]_A).$$

These two relations show that there is a correspondence between Lie algebroids structure on (A, M) and differentials on $\Gamma(\wedge A^*)$.

Theorem (Local splitting))

Let $x_0 \in M$ be a point where $\#_{x_0}$ has rank q . There exists a system of coordinates $(x_1, \dots, x_q, y_1, \dots, y_{n-q})$ valid in a neighborhood U of x_0 and a basis of sections $\{a_1, \dots, a_r\}$ of A over U , such that

$$\begin{aligned} \#(a_i) &= \partial_{x_i} \quad (i = 1, \dots, q), \\ \#(a_i) &= \sum_j b^{ij} \partial_{y_j} \quad (i = q+1, \dots, r), \end{aligned}$$

where $b^{ij} \in C^\infty(U)$ are smooth functions depending only on the y 's and vanishing at x_0 : $b^{ij} = b^{ij}(y^s)$, $b^{ij}(x_0) = 0$. Moreover, for any $i, j = 1, \dots, r$,

$$[a_i, a_j] = \sum_u C_{ij}^u a_u,$$

where $C_{ij}^u \in C^\infty(U)$ vanish if $u < q$ and satisfy $\sum_u \frac{\partial C_{ij}^u}{\partial y_j} b^{ut} = 0$.

Some examples of Lie algebroids

- (a) The basic example of a Lie algebroid over M is the tangent bundle itself, with the identity mapping as anchor.
- (b) Every finite dimensional Lie algebra is a Lie algebroid over a one point space.
- (c) Let (M, π) be a Poisson manifold. The bivector field π defines a bundle homomorphism $\pi_{\#} : T^*M \rightarrow TM$ and a bracket on $\Omega^1(M)$ by

$$[\alpha, \beta]_{\pi} = \mathcal{L}_{\pi_{\#}(\alpha)}\beta - \mathcal{L}_{\pi_{\#}(\beta)}\alpha - d\pi(\alpha, \beta)$$

such that $(T^*M, M, \pi_{\#}, [,]_{\pi})$ is a Lie algebroid.

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(d) Let $\mathfrak{g} \xrightarrow{\tau} \mathcal{X}(M)$ be an action of a finite-dimensional real Lie algebra \mathfrak{g} on a smooth manifold M , i.e., a morphism of Lie algebras from \mathfrak{g} to the Lie algebra of vector fields on M .

Consider $(A, M, \rho, [\cdot, \cdot]_A)$, where $A = M \times \mathfrak{g}$ as a trivial bundle and

$$\rho((m, \xi)) = \tau(\xi)(m) \quad \text{and} \quad [\xi, \eta]_A = \mathcal{L}_{\rho(\xi)}\eta - \mathcal{L}_{\rho(\eta)}\xi + [\xi, \eta]_{\mathfrak{g}},$$

$$\eta, \xi \in \Gamma(A) = C^\infty(M, \mathfrak{g}).$$

By using (ii) of Proposition 1.3, it is easy to check that $(A, M, \rho, [\cdot, \cdot]_A)$ is a Lie algebroid.

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By using (ii) of Proposition 1.3, it is easy to check that $(A, M, \rho, [,]_A)$ is a Lie algebroid.

Let $p : A \rightarrow M$ be a Lie algebroid with anchor map $\#$. An A -connection on a vector bundle $E \rightarrow M$ is an operator

$\nabla : \Gamma(A) \times \Gamma(E) \rightarrow \Gamma(E)$ satisfying:

- 1 $\nabla_{a+b}s = \nabla_a s + \nabla_b s$ for any $a, b \in \Gamma(A)$ and $s \in \Gamma(E)$;
- 2 $\nabla_a(s_1 + s_2) = \nabla_a s_1 + \nabla_a s_2$ for any $a \in \Gamma(A)$ and $s_1, s_2 \in \Gamma(E)$;
- 3 $\nabla_{fa}s = f \nabla_a s$ for any $a \in \Gamma(A)$, $s \in \Gamma(E)$ and $f \in C^\infty(M)$;
- 4 $\nabla_a(fs) = f \nabla_a s + \#(a)(f)s$ for any $a \in \Gamma(A)$, $s \in \Gamma(E)$ and $f \in C^\infty(M)$.

Definition

Let $p : A \rightarrow M$ be a Lie algebroid with anchor $\#$.

- ① An A -path is a smooth path $\alpha : [t_0, t_1] \rightarrow A$ such that

$$\#(\alpha(t)) = \frac{d}{dt}p(\alpha(t)), \quad t \in [t_0, t_1].$$

We call the curve $\gamma : [t_0, t_1] \rightarrow M$ given by $\gamma(t) = p(\alpha(t))$ the *base path* of α .

- ② An A -path α is called *vertical* if $\#(\alpha(t)) = 0$ for any $t \in [t_0, t_1]$.

Remark

When $A = TM$, a A -path is just the derivative $\dot{c} : [t_0, t_1] \rightarrow TM$ of a curve $c : [t_0, t_1] \rightarrow M$.

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Let $p : A \rightarrow M$ be a Lie algebroid, $\pi : E \rightarrow M$ a vector bundle and ∇ an A -connection on E . Fix an A -path $\alpha : [t_0, t_1] \rightarrow A$.

An α -section of E is a smooth map $s : [t_0, t_1] \rightarrow E$ such that the projections on M of α and s define the same base path, i.e.,

$$p(\alpha) = \pi(s).$$

We denote by $\Gamma(E)_\alpha$ the space of α -sections of E .

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Then there is exists an unique map

$$\nabla^\alpha : \Gamma(E)_\alpha \longrightarrow \Gamma(E)_\alpha$$

satisfying:

- 1 $\nabla^\alpha(c_1 s_1 + c_2 s_2) = c_1 \nabla^\alpha s_1 + c_2 \nabla^\alpha s_2$, $c_1, c_2 \in \mathbb{R}$;
- 2 $\nabla^\alpha f s = f' s + f \nabla^\alpha s$ where $f : [t_0, t_1] \longrightarrow \mathbb{R}$ is a smooth function;
- 3 if \tilde{s} is a local section of E which extends s then

$$\nabla^\alpha s(t) = \nabla_{\alpha(t)} \tilde{s};$$

Let (x_1, \dots, x_n) be a local system of coordinates on an open set U , (a_1, \dots, a_r) is local frame of A and (e_1, \dots, e_q) is a local frame of E over U and

$$\#a_k = \sum_{i=1}^n b^{ki} \partial_{x_i} \quad (k = 1, \dots, r).$$

Then

$$\alpha(t) = \sum_{i=1}^r \alpha_i(t) a_i \quad \text{and} \quad s(t) = \sum_{i=1}^q s_i(t) e_i.$$

We have

$$p(\alpha(t)) = \pi(s(t)) = (x_1(t), \dots, x_n(t)) \quad \text{and} \quad \#(\alpha(t)) = \sum_{i=1}^n x_i'(t) \partial_{x_i}.$$

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Then

$$\alpha(t) = \sum_{i=1}^r \alpha_i(t) a_i \quad \text{and} \quad s(t) = \sum_{i=1}^q s_i(t) e_i.$$

We have

$$p(\alpha(t)) = \pi(s(t)) = (x_1(t), \dots, x_n(t)) \quad \text{and} \quad \#(\alpha(t)) = \sum_{i=1}^n x_i'(t) \partial_{x_i}.$$

$$\begin{aligned}
 \nabla^\alpha s &= \sum_{i=1}^q s'_i(t) e_i + \sum_{i=1}^q s_i(t) \nabla^\alpha e_i \\
 &= \sum_{i=1}^q s'_i(t) e_i + \sum_{i=1}^q s_i(t) \nabla_{\alpha(t)} e_i \\
 &= \sum_{i=1}^q s'_i(t) e_i + \sum_{i=1}^q \sum_{j=1}^r s_i(t) \alpha_j(t) \nabla_{a_j} e_i \\
 &= \sum_{k=1}^q \left(s'_k(t) + \sum_{i=1}^q \sum_{j=1}^r s_i(t) \alpha_j(t) \Gamma_{ji}^k(x(t)) \right) e_k. \\
 \nabla_{a_j} e_i &= \sum_{k=1}^q \Gamma_{ji}^k e_k.
 \end{aligned}$$

An α -section s is called parallel along α if $\nabla^\alpha s = 0$, i.e.,

$$s'_k(t) + \sum_{i=1}^q \sum_{j=1}^r s_i(t) \alpha_j(t) \Gamma_{ji}^k(x(t)) = 0, \quad k = 1, \dots, q.$$

One has then the notion of parallel transport along α , denoted by

$$\tau_\alpha^t : E_{\gamma(t_0)} \longrightarrow E_{\gamma(t)},$$

and $\tau_\alpha^t(s_0) = s(t)$ where s is the unique parallel α -section satisfying $s(0) = s_0$.

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If $\alpha_0 \in A_x$ and s is a section of E in a neighborhood of x , one can check easily that

$$\nabla_{\alpha_0} s = \frac{d}{dt} \Big|_{t=0} (\tau_{\alpha}^t)^{-1}(s(\gamma(t))), \quad (4)$$

where α is any A -path satisfying $\alpha(0) = \alpha_0$.

Let $p : A \rightarrow M$ be a Lie algebroid with anchor $\#$. We shall call A -connections on the vector bundle $A \rightarrow M$ *linear A -connections*. Let \mathcal{D} be a linear A -connection. An A -path $\alpha : [t_0, t_1] \rightarrow A$ is a *geodesic* of \mathcal{D} if

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Let (x_1, \dots, x_n) be a local system of coordinates on an open set U and (a_1, \dots, a_r) a basis of local sections over U . The structure functions $b^{si}, C_{st}^u \in C^\infty(U)$ are given by

$$\#a_s = \sum_{i=1}^n b^{si} \partial_{x_i} \quad (s = 1, \dots, r),$$

We define the Christoffel symbols of \mathcal{D} according to (a_1, \dots, a_r) as usually by

$$\mathcal{D}_{a_s} a_t = \sum_{u=1}^r \Gamma_{st}^u a_u.$$

The A -path α is a geodesic if, for $i = 1, \dots, n$ and $j = 1, \dots, r$,

$$\left\{ \begin{array}{l} \dot{x}_i(t) = \sum_{j=1}^r \alpha_j(t) b^{ji}(x_1(t), \dots, x_n(t)), \\ \dot{\alpha}_j(t) = - \sum_{s,u=1}^r \alpha_s(t) \alpha_u(t) \Gamma_{su}^j(x_1(t), \dots, x_n(t)), \end{array} \right. \quad (5)$$

where $\alpha(t) = \sum_{i=1}^r \alpha_i(t) a_i$ is the local expression of α and $p(\alpha(t)) = (x_1(t), \dots, x_n(t))$ is the local expression of its base path.

Exactly as in the classical case, one has existence and uniqueness of geodesics with given initial base point $x \in M$ and "initial speed" $a_0 \in A_x$. Actually, there exists a vector field G on A such that the geodesics of \mathcal{D} are the integral curves of G . We call G the *geodesic vector field* associated to \mathcal{D} and \mathcal{D} is called complete if G is complete.

Remark

The notions of connection, parallel transport and geodesic can be defined in any anchored vector bundle.

We introduce now two natural notions of compatibility between linear A -connections and the structures of Lie algebroids.

Definition

- 1 *A linear A -connection \mathcal{D} is strongly compatible with the Lie algebroid structure if, for any A -path α , the parallel transport τ_α preserves $\ker \#$.*
- 2 *A linear A -connection \mathcal{D} is weakly compatible with the Lie algebroid structure if, for any vertical A -path α , the parallel transport τ_α preserves $\ker \#$.*

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The following proposition gives an useful characterization of the these notions of compatibility.

Proposition

- 1 *A linear A -connection \mathcal{D} is strongly compatible with the Lie algebroid structure if and only if, for any leaf L and for any sections $\alpha \in \Gamma(A_L)$ and $\beta \in \Gamma(\mathfrak{g}_L)$, $\mathcal{D}_\alpha \beta \in \Gamma(\mathfrak{g}_L)$.*
- 2 *A linear A -connection \mathcal{D} is weakly compatible with the Lie algebroid structure if and only if, for any leaf L and for any sections $\alpha \in \Gamma(\mathfrak{g}_L)$ and $\beta \in \Gamma(\mathfrak{g}_L)$, $\mathcal{D}_\alpha \beta \in \Gamma(\mathfrak{g}_L)$.*

Example

Let $p : A \rightarrow M$ be a Lie algebroid and ∇ be a TM-connection on A . Associated with ∇ there is an obvious linear A -connection

$$\mathcal{D}_a^0 b = \nabla_{\#(a)} b$$

which is clearly weakly compatible with the Lie algebroid structure.

A bit more subtle is the following linear A -connection

$$\mathcal{D}_a^1 b = \nabla_{\#(b)} a + [a, b]_A$$

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We give an interpretation of the torsion and the curvature of an A -connection which leads naturally to the notion of homotopy of A -paths. This notion plays a crucial role in the integrability of Lie algebroids.

Let $p : A \rightarrow M$ be a Lie algebroid with anchor $\#$ and $E \rightarrow M$ a vector bundle. The curvature of an A -connection ∇ on E is formally identical to the usual definition

$$R(a, b)s = \nabla_a \nabla_b s - \nabla_b \nabla_a s - \nabla_{[a, b]_A} s,$$

where $a, b \in \Gamma(A)$ and $s \in \Gamma(E)$. The connection ∇ is called flat if R vanishes identically.

If \mathcal{D} is a linear A -connection the torsion of \mathcal{D} is given by

$$T_{\mathcal{D}}(a, b) = \mathcal{D}_a b - \mathcal{D}_b a - [a, b]_A.$$

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In the classical case ($A = TM$), the curvature and the torsion can be interpreted by using variations of paths. We will show now that we have a similar interpretation in the general case. First, let us recall the notion of variation of paths in the classical case in order to find the appropriate generalization.

Let M be a manifold and ∇ a connection on TM .

A variation of curves is a smooth map $\Gamma = (-\varepsilon, \varepsilon) \times [a, b] \rightarrow M$. Any variation of curves defines two collections of curves: the **main curves** $\Gamma_s(t) = \Gamma(s, t)$ defined on $[a, b]$ by setting $s = \text{constant}$ and the **transverse curves** $\Gamma^{(t)}(s) = \Gamma(s, t)$ defined on $(-\varepsilon, \varepsilon)$ by setting $t = \text{constant}$.

The tangent vectors to these two families of curves are examples of vector fields along Γ , we denote them by

$$T(s, t) = \partial_t \Gamma(s, t) = \frac{d}{dt} \Gamma_s(t), \quad S(s, t) = \partial_s \Gamma(s, t) = \frac{d}{ds} \Gamma^{(t)}(s).$$

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If V is a vector field along Γ , we can compute the covariant derivative of V either along the main curves or along the transverse curves, the resulting vector fields along Γ are denoted by $D_t V$ and $D_s V$. The following lemma is classical.

Lemma

With the notation above, we have

$$D_s T - D_t S = T^\nabla(S, T). \quad (6)$$

and

$$D_s D_t Y - D_t D_s Y = -R^\nabla(S, T)Y. \quad (7)$$

Let (A, M, ρ) a Lie algebroid. Let $\alpha : [0, 1] \times [0, 1] \rightarrow A$ and $\Gamma : [0, 1] \times [0, 1] \rightarrow M$ its projection.

We call α *variation of A -paths* if

- for any $s \in [0, 1]$, the map $t \mapsto \alpha(s, t)$ is an A -path, i.e.,

$$\#(\alpha(s, t)) = \frac{\partial \Gamma}{\partial t}(s, t),$$

- the base variation $\Gamma(s, t) = \rho(\alpha(s, t))$ lies entirely in a fixed leaf L of the characteristic foliation.

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Let α be a variation of A -paths and Γ its projection.

A *transverse variation* to α is a smooth map
 $\beta : [0, 1] \times [0, 1] \longrightarrow A$ such that $p(\beta) = \Gamma$ and

$$\#(\beta(s, t)) = \frac{\partial \Gamma}{\partial s}(s, t).$$

It is clear that if $\#$ is injective, there is an unique transverse variation to a given variation of A -paths. However, if $\#$ is not injective, a given variation of A -paths admits many transverse variations to it. There is a way which permit the control of transverse variations to a fixed variation of A -path. Let us explain this important fact which is at the origin of the notion of homotopy of A -paths.

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First, let us fix some notations. Let α and β be, respectively, a variation of A -paths and a transverse variation and let $\Gamma = p(\alpha) = p(\beta)$ denote the commune base path.

Let ∇ be an A -connection on a vector bundle $\pi : E \rightarrow M$ and let $s : [0, 1] \times [0, 1] \rightarrow E$ be a section over Γ .

For any $\varepsilon \in [0, 1]$, $t \mapsto \alpha(\varepsilon, t)$ is an A -path and $\nabla_t s$ denotes the derivative of $t \mapsto s(\varepsilon, t)$ along this A -path.

On the other hand, for any $t \in [0, 1]$, $\varepsilon \mapsto \beta(\varepsilon, t)$ is an A -path and $\nabla_\varepsilon s$ denotes the derivative of $\varepsilon \mapsto s(\varepsilon, t)$ along this A -path.

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Proposition

With the notation above the following assertions hold.

- ① For any linear A -connection \mathcal{D} , the variation

$$\Delta(\alpha, \beta) = \mathcal{D}_t \beta - \mathcal{D}_\varepsilon \alpha - T_{\mathcal{D}}(\alpha, \beta)$$

does not depend on \mathcal{D} and satisfies $\#(\Delta(\alpha, \beta)) = 0$.

- ② for any A -connection ∇ on E and for any section s of E over Γ

$$\nabla_t \nabla_\varepsilon s - \nabla_\varepsilon \nabla_t s = R(\alpha, \beta)s + \nabla_{\Delta(\alpha, \beta)} s.$$

Proof.

Fix $(\varepsilon_0, t_0) \in [0, 1] \times [0, 1]$ and choose a local coordinates $(x_1, \dots, x_q, y_1, \dots, y_{n-q})$ near $x_0 = \Gamma(\varepsilon_0, t_0)$ and a basis of sections (a_1, \dots, a_r) , ($q = \text{rank} \#_{x_0}$), such that

$$\#(a_i) = \partial_{x_i} \quad (i = 1, \dots, q),$$

$$\#(a_i) = \sum_j b^{ij} \partial_{y_j} \quad (i = q + 1, \dots, r),$$

where $b^{ij} \in C^\infty(U)$ are smooth functions depending only on the y 's and vanishing at x_0 : $b^{ij} = b^{ij}(y^s)$, $b^{ij}(x_0) = 0$. Moreover, for any $i, j = 1, \dots, r$,

$$[a_i, a_j] = \sum_u C_{ij}^u a_u,$$

where $C_{ij}^u \in C^\infty(U)$ vanish if $u < q$ and satisfy $\sum_u \frac{\partial C_{ij}^u}{\partial y_j} b^{ut} = 0$.

In these coordinates, we have

$$\left\{ \begin{array}{l} \alpha(\varepsilon, t) = \sum_{i=1}^r \alpha^i(\varepsilon, t) a_i, \\ \beta(\varepsilon, t) = \sum_{i=1}^r \beta^i(\varepsilon, t) a_i, \\ \Gamma(\varepsilon, t) = (x_1(\varepsilon, t), \dots, x_q(\varepsilon, t), c_1, \dots, c_{n-q}), \\ \frac{\partial \Gamma}{\partial t} = \sum_{j=1}^q \frac{\partial x_j}{\partial t} \partial_{x_j} = \sum_{i=1}^q \alpha^i(\varepsilon, t) \partial_{x_i}, \\ \frac{\partial \Gamma}{\partial \varepsilon} = \sum_{j=1}^q \frac{\partial x_j}{\partial \varepsilon} \partial_{x_j} = \sum_{i=1}^q \beta^i(\varepsilon, t) \partial_{x_i}, \end{array} \right. \quad (8)$$

where c_1, \dots, c_{n-q} are constant.

Now

$$\mathcal{D}_t \beta = \sum_{i=1}^r \frac{\partial \beta^i}{\partial t} a_i + \sum_{i,j=1}^r \alpha^j \beta^i \mathcal{D}_{a_j} a_i$$

$$\mathcal{D}_\varepsilon \alpha = \sum_{i=1}^r \frac{\partial \alpha^i}{\partial \varepsilon} a_i + \sum_{i,j=1}^r \alpha^i \beta^j \mathcal{D}_{a_j} a_i.$$

Hence

$$\mathcal{D}_t \beta - \mathcal{D}_\varepsilon \alpha = \sum_{i=1}^r \left(\frac{\partial \beta^i}{\partial t} - \frac{\partial \alpha^i}{\partial \varepsilon} \right) a_i + T_{\mathcal{D}}(\alpha, \beta) + \sum_{i,j=1}^r \alpha^i \beta^j [a_i, a_j].$$

Now, from (8), we have $\frac{\partial \beta^i}{\partial t} = \frac{\partial \alpha^i}{\partial \varepsilon}$ for any $i = 1, \dots, q$, so

$$\mathcal{D}_t \beta - \mathcal{D}_\varepsilon \alpha - T_{\mathcal{D}}(\alpha, \beta) = \sum_{i=q+1}^r \left(\frac{\partial \beta^i}{\partial t} - \frac{\partial \alpha^i}{\partial \varepsilon} \right) a_i + \sum_{i,j=1}^r \alpha^i \beta^j [a_i, a_j]. \quad (9)$$

One can see that the right hand of this equality lies in $\ker \#$ and does not depend on \mathcal{D} .

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One can see that the right hand of this equality lies in $\ker \#$ and does not depend on \mathcal{D} .

We choose a local trivialization

$(x_1, \dots, x_q, y_1, \dots, y_{n-q}, a_1, \dots, a_r)$ as above, we trivialize E near x_0 by a local basis of sections (e_1, \dots, e_μ) and put

$$s(\varepsilon, t) = \sum_{j=1}^{\mu} s^j(\varepsilon, t) e_j.$$

$$\nabla_t s = \sum_{j=1}^{\mu} \frac{\partial s^j}{\partial t} e_j + \sum_{i,j} \alpha^i s^j \nabla_{a_i} e_j.$$

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$$\begin{aligned} \nabla_{\varepsilon} \nabla_t s &= \sum_{j=1}^{\mu} \frac{\partial^2 s^j}{\partial \varepsilon \partial t} e_j + \sum_{i,j} \left(\beta^i \frac{\partial s^j}{\partial t} + \frac{\partial \alpha^i}{\partial \varepsilon} s^j + \alpha^i \frac{\partial s^j}{\partial \varepsilon} \right) \nabla_{a_i} e_j \\ &\quad + \sum_{i,j,k} \beta^k \alpha^i s^j \nabla_{a_k} \nabla_{a_i} e_j. \end{aligned}$$

$$\begin{aligned} \nabla_t \nabla_{\varepsilon} s &= \sum_{j=1}^{\mu} \frac{\partial^2 s^j}{\partial t \partial \varepsilon} e_j + \sum_{i,j} \left(\alpha^i \frac{\partial s^j}{\partial \varepsilon} + \frac{\partial \beta^i}{\partial t} s^j + \beta^i \frac{\partial s^j}{\partial t} \right) \nabla_{a_i} e_j \\ &\quad + \sum_{i,j,k} \alpha^k \beta^i s^j \nabla_{a_k} \nabla_{a_i} e_j. \end{aligned}$$

$$\begin{aligned} \nabla_t \nabla_{\varepsilon} s - \nabla_{\varepsilon} \nabla_t s - R(\alpha, \beta)s &= \sum_{i,j} \left(\frac{\partial \beta^i}{\partial t} - \frac{\partial \alpha^i}{\partial \varepsilon} \right) s^j \nabla_{a_i} e_j \\ &\quad + \sum_{i,j,k} \alpha^k \beta^i s^j \nabla_{[a_k, a_i]} e_j. \end{aligned}$$

From the expression of $\Delta(\alpha, \beta)$ given by (9) and from (8), we have

$$\Delta(\alpha, \beta) = 0 \Leftrightarrow \begin{cases} \frac{\partial \alpha^i}{\partial \varepsilon} - \frac{\partial \beta^i}{\partial t} = \sum_{l,k=1}^r \alpha^l \beta^k C_{lk}^i & i = q+1, \dots, r, \\ \alpha^j = \frac{\partial x_j}{\partial t}, \beta^j = \frac{\partial x_j}{\partial \varepsilon} & j = 1, \dots, q. \end{cases} \quad (10)$$

Now by using the standard results about linear differential systems one can deduce easily the following useful proposition.

Proposition

Let $p : A \rightarrow M$ be a Lie algebroid. Then, for a given variation of A -paths α and for given $\beta_0 : [0, 1] \rightarrow A$ such that

$$\#(\beta_0)(\varepsilon) = \frac{\partial p \circ \alpha}{\partial \varepsilon}(\varepsilon, 0)$$

there exists an unique transverse variation β to α such that

$$\Delta(\alpha, \beta) = 0 \quad \text{and} \quad \beta(\varepsilon, 0) = \beta_0(\varepsilon) \quad \text{for any } \varepsilon \in [0, 1].$$

We can now define the homotopy of A -paths with fixed end-points.

Definition

Let α_0 and α_1 be two A -paths on a Lie algebroid $p : A \rightarrow M$ such that $p(\alpha_0(0)) = p(\alpha_1(0))$ and $p(\alpha_0(1)) = p(\alpha_1(1))$.

An A -homotopy with fixed end-points from α_0 to α_1 is a variation of A -paths α such that:

- 1 $p(\alpha(\varepsilon, 0)) = p(\alpha(0, 0))$ and $p(\alpha(\varepsilon, 1)) = p(\alpha(0, 1))$ for any $\varepsilon \in [0, 1]$, $\alpha(0, \cdot) = \alpha_0$ and $\alpha(1, \cdot) = \alpha_1$,
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The following Lemma will be useful latter.

Lemma

Let $\alpha_0 : [0, 1] \rightarrow A$ be an A -path and $\beta_0 : [0, 1] \rightarrow A$ an α_0 -section such that $\beta_0(0) = \beta_0(1) = 0$. Then there exists an A -homotopy α with fixed end-points such that $\alpha(0, \cdot) = \alpha_0$ and the corresponding transverse variation β satisfies $\beta(0, \cdot) = \beta_0$.

Proof.

Consider the base path $\gamma_0 : [0, 1] \rightarrow M$ of α_0 and choose an homotopy $\gamma : [0, 1] \times [0, 1] \rightarrow M$ with fixed end points such that γ lies in the same leaf as γ_0 , $\gamma(0, \cdot) = \gamma_0$ and $\frac{\partial \gamma}{\partial \varepsilon}(0, t) = \#(\beta_0(t))$. We choose also $\beta : [0, 1] \times [0, 1] \rightarrow A$ such that $\beta(0, t) = \beta_0(t)$ for any $t \in [0, 1]$, $\beta(\varepsilon, 0) = \beta(\varepsilon, 1) = 0$ for any $\varepsilon \in [0, 1]$ and $\frac{\partial \gamma}{\partial \varepsilon}(\varepsilon, t) = \#(\beta(\varepsilon, t))$ for any (ε, t) . From (10), one can deduce that there exists an unique variation $\alpha : [0, 1] \times [0, 1] \rightarrow A$ such that the base path of α is γ , $\frac{\partial \alpha}{\partial t}(\varepsilon, t) = \#(\alpha(\varepsilon, t))$, $\alpha(0, \cdot) = \alpha_0$ and $\Delta(\alpha, \beta) = 0$. This variation is clearly an A -homotopy with fixed end-points and satisfies the required properties.

A Riemannian metric on a Lie algebroid $p : A \rightarrow M$ is the data, for any $x \in M$, of a scalar product $\langle \cdot, \cdot \rangle_x$ on the fiber A_x such that, for any local sections $a, b \in \Gamma(A)$, the function $\langle a, b \rangle$ is smooth.

the formula

$$2\langle \mathcal{D}_a b, c \rangle = \#(a) \cdot \langle b, c \rangle + \#(b) \cdot \langle a, c \rangle - \#(c) \cdot \langle a, b \rangle \\ + \langle [c, a], b \rangle + \langle [c, b], a \rangle + \langle [a, b], c \rangle$$

defines a linear A -connection which is characterized by the two following properties:

- (i) \mathcal{D} is metric, i.e., $\#(a) \cdot \langle b, c \rangle = \langle \mathcal{D}_a b, c \rangle + \langle b, \mathcal{D}_a c \rangle$,
- (ii) \mathcal{D} is torsion free, i.e., $\mathcal{D}_a b - \mathcal{D}_b a = [a, b]$.

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In a system of coordinates (x_1, \dots, x_n) over a trivializing neighborhood U of M where A admits a basis of local sections (a_1, \dots, a_r) the Levi-Civita A -connection is determined by the Christoffel's symbols defined by $\mathcal{D}_{a_i} a_j = \sum_k \Gamma_{ij}^k a_k$. A direct computation gives

$$\begin{aligned} \Gamma_{ij}^k &= \frac{1}{2} \sum_{l=1}^r \sum_{u=1}^n g^{kl} \left(b^{iu} \partial_{x_u} (g_{jl}) + b^{ju} \partial_{x_u} (g_{il}) - b^{lu} \partial_{x_u} (g_{ij}) \right) \\ &\quad + \frac{1}{2} \sum_{l=1}^r \sum_{u=1}^r g^{kl} (C_{ij}^u g_{ul} + C_{li}^u g_{uj} + C_{lj}^u g_{ui}), \end{aligned} \quad (11)$$

where the structure functions $b^{si}, C_{st}^u \in C^\infty(U)$ are given by

$$\#a_s = \sum_{i=1}^n b^{si} \partial_{x_i} \quad (s = 1, \dots, r),$$

Remark

There are two extremal cases:

- 1 The Lie algebroid A is the tangent bundle TM of a manifold and we recover the classical notion of Riemannian manifold.
- 2 The Lie algebroid A is a Lie algebra \mathfrak{g} considered as a Lie algebroid over a point. In this case a Riemannian metric on \mathfrak{g} is a scalar product $\langle \cdot, \cdot \rangle$ and the Levi-Civita \mathfrak{g} -connection is the product $\mathcal{D} : \mathfrak{g} \times \mathfrak{g} \rightarrow \mathfrak{g}$ given by

$$2\langle \mathcal{D}_u v, w \rangle = \langle [u, v], w \rangle + \langle [w, u], v \rangle + \langle [w, v], u \rangle.$$

Actually \mathcal{D} is the infinitesimal data associated to the Levi-Civita connection of the left invariant metric associated to $\langle \cdot, \cdot \rangle$ on any Lie group with \mathfrak{g} as a Lie algebra.

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Let $\langle \cdot, \cdot \rangle$ be a Riemannian metric on a Lie algebroid $p : A \rightarrow M$ with anchor $\#$. For any leaf L of the characteristic foliation and for any $x \in L$,

$$A_x = \mathfrak{g}_x \oplus \mathfrak{g}_x^\perp,$$

where \mathfrak{g}_x^\perp is the orthogonal to \mathfrak{g}_x with respect $\langle \cdot, \cdot \rangle_x$. The restriction of the anchor $\#$ to \mathfrak{g}_x^\perp is an isomorphism into $T_x L$ and hence induces a scalar product on $T_x L$

$$\langle u, v \rangle_L = \langle a, b \rangle,$$

where $a, b \in \mathfrak{g}_x^\perp$ and $\#(a) = u$ and $\#(b) = v$. Thus $\langle \cdot, \cdot \rangle$ induces a Riemannian metric $\langle \cdot, \cdot \rangle_L$ on L . We call it the *induced Riemannian metric* on L . On the other hand, the scalar product $\langle \cdot, \cdot \rangle_x$ induces a scalar product on \mathfrak{g}_x and we denote by \widehat{D} the Levi-Civita \mathfrak{g}_x -connection associated with $(\mathfrak{g}_x, \langle \cdot, \cdot \rangle_x)$.

Let us precise more this situation. Fix a leaf L and consider $\rho_L : A_L \longrightarrow L$. We have

$$A_L = \mathfrak{g}_L \oplus \mathfrak{g}_L^\perp.$$

We call the elements of $\Gamma(\mathfrak{g}_L)$ *vertical sections* and the elements of $\Gamma(\mathfrak{g}_L^\perp)$ *horizontal sections*. For any section a , we denote by a^v and a^h , respectively, its horizontal and vertical component. Note that the bracket of a vertical section with every section is a vertical section. Thus, in the Riemannian point of view, the short exact sequence

$$0 \longrightarrow \mathfrak{g}_L \longrightarrow A_L \longrightarrow TL$$

is formally identical to a Riemannian submersion. So we can introduce the O'Neill tensors.

We denote by T and H the elements of $\Gamma(A^* \otimes A^* \otimes A)$ whose values on sections a, b are given by

$$T_a b = (\mathcal{D}_{a^\vee} b^\vee)^h + (\mathcal{D}_{a^\vee} b^h)^\vee \quad \text{and} \quad H_a b = (\mathcal{D}_{a^h} b^\vee)^h + (\mathcal{D}_{a^h} b^h)^\vee.$$

The following properties of T and H follow immediately from the definition: for any $a, b \in \Gamma(A)$,

$$H_{a^h} b^h = \frac{1}{2}[a^h, b^h]^\vee, \tag{12}$$

$$\mathcal{D}_{a^\vee} b^h = T_{a^\vee} b^h + (\mathcal{D}_{a^\vee} b^h)^h, \tag{13}$$

$$\mathcal{D}_{a^h} b^\vee = (\mathcal{D}_{a^h} b^\vee)^\vee + H_{a^h} b^\vee, \tag{14}$$

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Moreover, for any $u, v \in \mathfrak{g}_x$,

$$\mathcal{D}_u v = \widehat{\mathcal{D}}_u v + T_u v. \quad (16)$$

The following proposition is an immediate consequence of (15).

Proposition

Let $\gamma : [t_0, t_1] \rightarrow L$ be a smooth path and let $\gamma^h : [t_0, t_1] \rightarrow \mathfrak{g}_L^\perp$ be the unique A -path with the base path γ . Then γ is a geodesic with respect to the induced Riemannian metric on L if and only if γ^h is a geodesic of the Levi-Civita A -connexion.

The following proposition gives an interpretation of the tensors T and H .

Proposition

- 1 *The Levi-Civita A -connection is strongly compatible with the Lie algebroid structure if and only if $T = H = 0$.*
- 2 *The Levi-Civita A -connection is weakly compatible with the Lie algebroid structure if and only if $T = 0$.*

Let $p : A \longrightarrow M$ be a Lie algebroid and $\langle \cdot, \cdot \rangle$ a Riemannian metric on A . The Riemannian metric defines a bundle isomorphism between A and A^* which transport the Lie-Poisson structure on A^* into a Poisson structure say $\pi_{\langle \cdot, \cdot \rangle}$ in A . Let $E : A \longrightarrow \mathbb{R}$ be the energy function given by $E(a) = \frac{1}{2}\langle a, a \rangle$ and let X_E denote the Hamiltonian vector field associated to E with respect to $\pi_{\langle \cdot, \cdot \rangle}$. The following result is a generalization of a well-known result in Riemannian geometry.

Theorem

The geodesics of the Levi-Civita A -connection associated to $\langle \cdot, \cdot \rangle$ are the integral curves of the Hamiltonian vector field X_E .

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The flow of the Hamiltonian vector field X_E is called the *geodesic flow* of $\langle \cdot, \cdot \rangle$.

Remark

Let $p : A \rightarrow M$ be a Riemannian Lie algebroid. Then:

- 1 For any leaf L , the geodesic vector field X_E is tangent to A_L and to \mathfrak{g}_x for any $x \in L$.
- 2 From Proposition 2.1, one can deduce that, for any leaf L , the geodesic vector field X_E is tangent to \mathfrak{g}_L^\perp .

Corollary

Let $p : A \rightarrow M$ be Riemannian Lie algebroid. Then

- 1 If L is a compact leaf then the geodesic flow is complete in restriction to A_L .
- 2 If M is compact then the geodesic flow is complete and for any leaf L , the induced Riemannian metric $\langle \cdot, \cdot \rangle_L$ is complete.

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Sasaki Metric of a Riemannian Lie algebroid

Let $p : A \rightarrow M$ be a Riemannian Lie algebroid with anchor $\#$. Fix a leaf L , consider $p_L : A_L \rightarrow L$ and put $\mathcal{V}A_L = \text{Ker} dp_L$.

For any $a \in A_L$, we consider the subspace $\mathfrak{h}^\perp A_L$ of $T_a A_L$ consisting of the tangent vectors V_a such that there exists an horizontal A -path $\alpha : [0, 1] \rightarrow \mathfrak{g}_L^\perp$ satisfying $p(\alpha(0)) = p(a)$ and

$V_a = \frac{d}{dt}\bigg|_{t=0} \tau_\alpha^t(a)$, where τ_α is the parallel transport along α . We have

$$TA_L = \mathcal{V}A_L \oplus \mathfrak{h}^\perp A_L. \quad (17)$$

Indeed, we define $K : TA_L \longrightarrow A_L$ as follows. Fix $a \in A_L$ and $Z \in T_a A_L$ and choose $\beta : [0, 1] \longrightarrow A_L$ such that $\beta(0) = a$ and $\dot{\beta}(0) = Z$. There exists a unique horizontal A -path $\alpha : [0, 1] \longrightarrow \mathfrak{g}_L^\perp$ with the base path $p \circ \beta(t)$. Put

$$K(Z) = (\mathcal{D}^\alpha \beta)(0).$$

It is easy to check that K is well-defined, $\text{Ker} K = \mathfrak{h}^\perp A_L$ and, for any $Z \in \mathcal{V}A_L$, $K(Z) = Z$. Then the relation (17) follows.

Let (x_1, \dots, x_l) be a system of local coordinates on an open set U in L and (a_1, \dots, a_r) is a basis of local sections (over U) of A_L . This defines a system of coordinates $(x_1, \dots, x_l, \mu_1, \dots, \mu_r)$ on A_L and if $Z = \sum_j b_j \partial_{x_j} + \sum_j Z^j \partial_{\mu_j}$ then

$$K(Z) = \sum_l \left(Z^l + \sum_{i,j} \alpha_i \mu_j \Gamma_{ij}^l \right) a_l, \quad (18)$$

where $dp_L(Z) = \#(\sum_i \alpha_i a_i)$ and $\sum_i \alpha_i a_i \in \mathfrak{g}_L^\perp$.

Remark

In general, the geodesic vector field does not lie in $\text{Ker}K$. Indeed, one can check easily that for any $a \in A_L$

$$K(X_E(a)) = -\mathcal{D}_{a^\vee} a.$$

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We define the Sasaki metric on A_L by

$$g_L(Z_a, Z_a) = \langle d_a p(Z_a), d_a p(Z_a) \rangle_L + \langle K(Z_a), K(Z_a) \rangle.$$

The projection $p_L : A_L \rightarrow L$ becomes a Riemannian submersion. We consider now the Liouville vector field \vec{r} on A_L which is the vector field generating the flow $\phi_t(a) = e^t a$. By direct computation one can get

$$[\vec{r}, X_E] = X_E. \quad (19)$$

From this relation, one deduce that X_E preserves the Riemannian volume on A_L associated to g_L if and only if X_E preserves the Riemannian volume of the restriction of g_L to the spheres bundle $UA_L = \{a \in A_L; \langle a, a \rangle = 1\}$. Let us compute the divergence of the geodesic vector field with respect to g_L .

Theorem

The divergence the geodesic vector field X_E with respect to the Sasaki metric g_L is given by

$$\operatorname{div}(X_E)(a) = \operatorname{Tr} ad_{a^\vee} + \langle a^h, N \rangle, \quad (20)$$

where $ad_{a^\vee} : \mathfrak{g}_{p(a)} \longrightarrow \mathfrak{g}_{p(a)}$, $b \longrightarrow [a^\vee, b]$ and $N = \sum_i T_{b_i} b_i$ where (b_1, \dots, b_s) is any orthonormal basis of $\mathfrak{g}_{p(a)}$ and T is the O'Neill tensor.

The following proposition gives an interesting interpretation of $\operatorname{div}X_E$, namely $\operatorname{div}X_E$ is a modular cocycle.

Proposition

*Let $p : A \rightarrow M$ be a transitive Riemannian Lie algebroid such that both A and TM are orientable. Denote by $\lambda \in \Gamma(\wedge^{\operatorname{top}} A)$ and $\nu \in \Gamma(\wedge^{\operatorname{top}} T^*M)$, respectively, the Riemannian volume associated to \langle , \rangle and the Riemannian volume associated to \langle , \rangle_M then*

$$\mathcal{D}^A(\lambda \otimes \nu) = \operatorname{div}(X_E)(\lambda \otimes \nu),$$

where \mathcal{D}^A is the canonical representation of A . Thus $\operatorname{div}(X_E)$ is a modular cocycle.

Remark

- 1 If $A = TM$ then $\operatorname{div}(X_E) = 0$ and one recovers the classical Liouville Theorem.
- 2 If A is a Lie algebra then $\operatorname{div}(X_E) = 0$ if and only if A is unimodular.
- 3 If A is a transitive unimodular Lie algebroid then there exists a Riemannian metric on A such that $\operatorname{div}(X_E) = 0$.

Let $p : A \rightarrow M$ be a Riemannian Lie algebroid with anchor $\#$.
For any A -path $\alpha : [0, 1] \rightarrow A$, the energy and the length of α are given, respectively, by

$$\mathbf{E}(\alpha) = \frac{1}{2} \int_0^1 \langle \alpha(t), \alpha(t) \rangle dt \quad \text{and} \quad \mathcal{L}(\alpha) = \int_0^1 \sqrt{\langle \alpha(t), \alpha(t) \rangle} dt.$$

For any m, q lying in the same leaf of the characteristic foliation, we denote by Ω_{mq} the set of A -path α such that $p(\alpha(0)) = m$ and $p(\alpha(1)) = q$.

Proposition

(First variation formulas) Let $p : A \rightarrow M$ be a Riemannian Lie algebroid. Then:

- 1 For any variation of A -paths $\alpha : [0, 1] \times [0, 1] \rightarrow A$ and for any β a transverse variation to α , one has

$$\begin{aligned} \frac{d}{d\varepsilon} \mathbf{E}(\alpha) &= \langle \beta(\varepsilon, 1), \alpha(\varepsilon, 1) \rangle - \langle \beta(\varepsilon, 0), \alpha(\varepsilon, 0) \rangle - \int_0^1 \langle \beta, \mathcal{D}_t \alpha \rangle dt \\ &\quad - \int_0^1 \langle \Delta(\alpha, \beta), \alpha \rangle dt. \end{aligned}$$

- 2 The h -critical points of $\mathbf{E} : \Omega_{mq} \rightarrow \mathbb{R}$, namely the A -paths α_0 such that

$$\frac{d}{d\varepsilon} \mathbf{E}(\alpha)|_{\varepsilon=0} = 0$$

Proposition

(Second variation formulas) *Let $p : A \rightarrow M$ be a Riemannian Lie algebroid. Then the following assertions hold. For any variation of A -paths α such that α_0 is a geodesic and for any β a transverse variation to α such that $\Delta(\alpha, \beta) = 0$, one has*

$$\begin{aligned} \frac{d^2}{d\varepsilon^2} \mathbf{E}(\alpha)|_{\varepsilon=0} &= \langle \mathcal{D}_\varepsilon \beta(0, 1), \alpha(0, 1) \rangle - \langle \mathcal{D}_\varepsilon \beta(0, 0), \alpha(0, 0) \rangle \\ &+ \int_0^1 \langle \mathcal{D}_t \beta_0, \mathcal{D}_t \beta_0 \rangle dt + \int_0^1 \langle \beta_0, R(\alpha_0, \beta_0) \alpha_0 \rangle dt. \end{aligned}$$

Proposition

Let α be an A -homotopy of A -paths such that α_0 is a geodesic and let β be the corresponding transverse variation. One has

$$\frac{d^2}{d\varepsilon^2} \mathbf{E}(\alpha)|_{\varepsilon=0} = \int_0^1 \langle \mathcal{D}_t \beta_0, \mathcal{D}_t \beta_0 \rangle dt + \int_0^1 \langle \beta_0, R(\alpha_0, \beta_0) \alpha_0 \rangle dt.$$

Proposition

Let α be a variation of A -paths such that α_0 is a geodesic parameterized by arc length and let β a transverse variation to α such that $\Delta(\alpha, \beta) = 0$. One has

$$\begin{aligned} \frac{d^2}{d\varepsilon^2} \mathcal{L}(\alpha)|_{\varepsilon=0} &= \langle \mathcal{D}_\varepsilon \beta(0, 1), a(0, 1) \rangle - \langle \mathcal{D}_\varepsilon \beta(0, 0), \alpha(0, 0) \rangle \\ &+ \int_0^1 \langle \mathcal{D}_t \beta_0, \mathcal{D}_t \beta_0 \rangle dt + \int_0^1 \langle \beta_0, R(\alpha_0, \beta_0) \alpha_0 \rangle dt \\ &- \int_0^1 \langle \alpha_0, \mathcal{D}_t \beta_0 \rangle dt. \end{aligned}$$

Proposition

Let α be an A -homotopy of A -paths such that α_0 is a geodesic parameterized by arc length and let β be the corresponding transverse variation. One has

$$\begin{aligned} \frac{d^2}{d\varepsilon^2} \mathcal{L}(\alpha)|_{\varepsilon=0} &= \int_0^1 \langle \mathcal{D}_t \beta_0, \mathcal{D}_t \beta_0 \rangle dt + \int_0^1 \langle \beta_0, R(\alpha_0, \beta_0) \alpha_0 \rangle dt \\ &\quad - \int_0^1 \langle \alpha_0, \mathcal{D}_t \beta_0 \rangle dt. \end{aligned}$$

As an application of Proposition 4.1 2., we give now a description of the geodesics of a left invariant Riemannian metric on a Lie group using the geodesics of its Lie algebra considered as a Riemannian Lie algebroid.

Let G be a Lie group and $\mathfrak{g} = T_e G$ its Lie algebra. For any $u \in \mathfrak{g}$, we denote by u^+ the associated left invariant vector field on G . Suppose that G is endowed with a left invariant Riemannian metric g and put $\langle \cdot, \cdot \rangle = g_e$. If we think \mathfrak{g} as a Lie algebroid, $(\mathfrak{g}, \langle \cdot, \cdot \rangle)$ is a Riemannian Lie algebroid and we will explain how one can construct the geodesics of (G, g) from the geodesics of $(\mathfrak{g}, \langle \cdot, \cdot \rangle)$.

Choose a basis (e_1, \dots, e_n) of \mathfrak{g} and put $g_{ij} = \langle e_i, e_j \rangle$. Recall that the geodesics of $(\mathfrak{g}, \langle \cdot, \cdot \rangle)$ are the integral curves of the geodesic vector field X_E given in the linear coordinates (x_1, \dots, x_n) associated to (e_1, \dots, e_n) by

$$X_E = - \sum_{s,t,j} x_s x_t \Gamma_{st}^j \partial_{x_j},$$

where Γ_{st}^j are given by

$$\Gamma_{st}^j = \frac{1}{2} \sum_{l,u} g^{lj} (g_{ul} C_{st}^u + g_{ut} C_{ls}^u + g_{us} C_{lt}^u).$$

Here (g^{ij}) is the inverse matrix of (g_{ij}) and C_{ij}^k are given by $[e_i, e_j] = \sum_u C_{ij}^u e_u$.

Proposition

Let $h \in G$ and $v \in T_h G$. Then the geodesic $\gamma : \mathbb{R} \rightarrow G$ of (G, g) satisfying $\gamma(0) = h$ and $\dot{\gamma}(0) = v$ is the integral curve passing through h of the time-dependent family of left invariant vector fields $(\alpha^+(t))_{t \in \mathbb{R}}$ where $\alpha : \mathbb{R} \rightarrow \mathfrak{g}$ is the geodesic of $(\mathfrak{g}, \langle \cdot, \cdot \rangle)$ satisfying $\alpha(0) = (L_{h^{-1}})_*(v)$.

Remark

If the Riemannian metric g is bi-invariant then $\Gamma_{ij}^k = \frac{1}{2} C_{ij}^k$ and hence X_E vanishes identically. We deduce from Proposition 3.5 that the geodesic of (G, g) passing through $h \in G$ and with initial velocity $v \in T_h G$ is the integral curve (passing through h) of the left invariant vector field $(L_{h^{-1}*}(v))^+$.

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Let $p : A \rightarrow M$ be a Riemannian Lie algebroid. The different curvatures (sectional curvature, Ricci curvature and scalar curvature) can be defined as the classical case (when $A = TM$). For any leaf L , the short exact sequence

$$0 \rightarrow \mathfrak{g}_L \rightarrow A_L \rightarrow TL$$

is formally identical to a Riemannian submersion and hence all formulas on curvature given by O'Neill are valid in this context. We denote by K , \hat{K} and \tilde{K} respectively, the sectional curvature of the Riemannian metrics $\langle \cdot, \cdot \rangle$, the restriction of $\langle \cdot, \cdot \rangle$ to \mathfrak{g}_L and the induced metric on L . The following proposition is a reformulation of Corollary 9.29 pp. 241 in [?].

Proposition

Let $\alpha, \beta, s_1, s_2 \in \Gamma(A_L)$ such that α, β are vertical, s_1, s_2 are horizontal and $|\alpha \wedge \beta| = 1$, $|s_1| = |\alpha| = 1$, $|s_1 \wedge s_2| = 1$. Then

$$\begin{aligned}K(\alpha, \beta) &= \hat{K}(\alpha, \beta) + |T_\alpha \beta|^2 - \langle T_\alpha \alpha, T_\beta \beta \rangle, \\K(s_1, \alpha) &= \langle (\mathcal{D}_{s_1} T)_\alpha \alpha, s_1 \rangle - |T_\alpha s_1| + |H_{s_1} \alpha|^2, \\K(s_1, s_2) &= \tilde{K}(s_1, s_2) - 3|H_{s_1} s_2|^2.\end{aligned}$$

The last formula says that the leaves carry "more curvature" than the Lie algebroid and by applying Mayer theorem (see for instance [?]) we get:

Proposition

Let $A \rightarrow M$ be a complete Riemannian algebroid and let L be a leaf of the characteristic foliation such that for any linearly independent horizontal sections s_1, s_2 over L , $K(s_1, s_2) \geq k$. Then $\text{diam} L \leq \frac{\pi}{\sqrt{k}}$ and hence L is compact.

There is another case when one can apply Mayer theorem. Consider a Riemannian Lie algebroid $p : A \rightarrow M$ such that the O'Neill tensor T vanishes and fix a leaf L and denote by r and \tilde{r} respectively the Ricci curvature of the Riemannian metrics $\langle \cdot, \cdot \rangle$ and $\langle \cdot, \cdot \rangle_L$. The formula 9.36c pp.244 in [?] applies in our context and gives

$$r(s_1, s_2) = \tilde{r}(\#(s_1), \#(s_2)) - 2 \sum_{i=1}^l \langle H_{s_1} a_i, H_{s_2} a_i \rangle$$

where (a_1, \dots, a_l) is any orthonormal basis of \mathfrak{g}_L^\perp .

By applying Mayer theorem (see for instance [?]) we get:

Proposition

Let $A \rightarrow M$ be a complete Riemannian algebroid such that $T = 0$ and let L be a leaf of the characteristic foliation such that there exists a constant k such that the restriction of r to \mathfrak{g}_L^\perp satisfies

$$r \geq (n - 1)k^{-2} \langle \cdot, \cdot \rangle.$$

Then $\text{diam} L \leq \frac{\pi}{\sqrt{k}}$ and hence L is compact.

A groupoid is a small category \mathcal{C} in which all the arrows are invertible. We shall write M for the set of objects of \mathcal{C} , while the set of arrows of \mathcal{C} will be denoted by \mathcal{C} . We shall often identify M with the subset of units of \mathcal{C} . The structure maps of \mathcal{C} will be denoted as follows: $\mathbf{s}, \mathbf{t} : \mathcal{C} \rightarrow M$ will stand for the source map, respectively the target map, $m : \mathcal{C}^2 = \{(g, h); \mathbf{s}(g) = \mathbf{t}(h)\} \rightarrow \mathcal{C}$ the multiplication map ($m(g, h) = gh$), $i : \mathcal{C} \rightarrow \mathcal{C}_1$ ($i(g) = g^{-1}$) for the inverse map and $u : M \rightarrow \mathcal{C}$ ($u(x) = 1_x$) for the unit map. Given $g \in \mathcal{C}$, the right multiplication by g is only defined on the \mathbf{s} -fiber at $\mathbf{t}(g)$, and induces a bijection

$$R_g : \mathbf{s}^{-1}(\mathbf{t}(g)) \rightarrow \mathbf{s}^{-1}(\mathbf{s}(g)).$$

A Lie groupoid is a groupoid \mathcal{C} , equipped with the structure of smooth manifold both on the \mathcal{C} and on the M such that all the structure maps are smooth and \mathbf{s} and \mathbf{t} are submersions.

The construction of a Lie algebra of a given Lie group extends to Lie groupoids. Explicitly, if \mathcal{C} is a Lie groupoid, the vector bundle $T^{\mathbf{s}}\mathcal{C} = \text{Ker}(d\mathbf{s})$ over \mathcal{C} of \mathbf{s} -vertical tangent vectors pulls back along $i : M \rightarrow \mathcal{C}$ to a vector bundle A over M . This vector bundle has the structure of a Lie algebroid. Its anchor $\# : A \rightarrow TM$ is induced by the differential of the target map, $d\mathbf{t} : T\mathcal{C} \rightarrow TM$. The sections of A over M can be identified by the space of right invariant \mathbf{s} -vertical vector fields which induce a Lie bracket on the space of sections of A .

With this construction in mind, one can see that a Riemannian structure on A is equivalent to the data of a Riemannian metric on any \mathbf{s} -fiber such that, for any $g \in \mathcal{C}$, $R_g : \mathbf{s}^{-1}(\mathbf{t}(g)) \rightarrow \mathbf{s}^{-1}(\mathbf{s}(g))$ is an isometry. In this case, for any $x \in M$, $\mathbf{t} : \mathbf{s}^{-1}(x) \rightarrow L_x$ is a Riemannian submersion where the leaf L_x is endowed with the metric defined in 3.1.

A Lie algebroid A is called integrable if it is isomorphic to the Lie algebroid associated to a Lie groupoid. In [?], Crainic and Fernandes give a final solution to the problem of integrability of Lie algebroids. They show that the obstruction to integrability can be controlled by two computable quantities. The following proposition is a direct application of Crainic-Fernandes results on integrability.

Proposition

Let $p : A \rightarrow M$ be a Riemannian Lie algebroid such that $H = 0$. Then A is integrable.

There is a large class of Lie algebroids for which one can apply this result. Let (M, π) be a Poisson manifold. The cotangent bundle T^*M carries a structure of a Lie algebroid where the anchor is the contraction by π , $\pi_{\#} : T^*M \rightarrow TM$ and the Lie bracket is given by the Koszul bracket

$$[\alpha, \beta] = \mathcal{L}_{\pi_{\#}(\alpha)}\beta - \mathcal{L}_{\pi_{\#}(\beta)}\alpha - d\pi(\alpha, \beta)$$

where $\alpha, \beta \in \Omega^1(M)$. Let $\langle \cdot, \cdot \rangle$ be a Riemannian structure in T^*M . In [?], the author studied the triple $(M, \pi, \langle \cdot, \cdot \rangle)$ such that π is parallel with respect the Levi-Civita T^*M -connection \mathcal{D} . A triple $(M, \pi, \langle \cdot, \cdot \rangle)$ satisfying $\mathcal{D}\pi = 0$ is called Riemann-Poisson manifold. The condition $\mathcal{D}\pi = 0$ implies that $\text{Ker}\pi_{\#}$ is invariant by parallel transport and hence \mathcal{D} is strongly compatible with the Lie algebroid structure of T^*M . By using Proposition 2.2, we deduce that $H = 0$.

So we get the following result.

Corollary

*Let $(M, \pi, \langle \cdot, \cdot \rangle)$ be a Riemann-Poisson manifold. Then the Lie algebroid structure of T^*M associated to π is integrable.*